

Some Features of the Hadronic $B_c^{(*)}$ -meson Production at Large p_T

A.V. Berezhnoy, A.K. Likhoded, O.P. Yushchenko
Institute for High Energy Physics, Protvino 142284, Russia
E-mail: LIKHODED@MX.IHEP.SU

Calculations of the hadronic $B_c^{(*)}$ -mesons production performed in the framework of the perturbative QCD taking into account $O(\alpha_s^4)$ Feynmann diagrams are presented. A comparison of the exact calculations with those based on the fragmentation model of $\bar{b} \rightarrow B_c^{(*)} + X$ shows the large discrepancy between them. The exact calculations of the $B_c^{(*)}$ -mesons production cross-sections as the function of p_T at the energy of the FNAL Tevatron ($\sqrt{s} = 1.8$ TeV) are given. The predicted ratio of the vector to the pseudoscalar state cross-sections is about $R \sim 3$ instead of $R \sim 1.4$ for the fragmentation model.

I. INTRODUCTION

The production of hadrons with two heavy constituent quarks of different flavours was the subject of the theoretical studies over two the last years. The calculations of the $B_c^{(*)}$ -meson production cross-section in e^+e^- annihilation at the Z^0 pole [1-3] show a principal possibility of the ground state observation of the $\bar{b}c$ system as well as its higher excitations at LEP-I energies. It is much more possible to observe $B_c^{(*)}$ -meson in experiments at the hadronic collider where we can expect [4-6] a large number of events at the available luminosity of the FNAL collider.

Another aspect of interest to the processes of the heavy quark system production consists in the fact that it offers a possibility to understand the mechanism of the hadronization in the framework of the perturbative QCD with the minimal assumptions concerning wave functions of the bound states.

Let us enumerate the main results of the studies of the $\bar{b}c$ system production mechanisms in different reactions:

i) $B_c^{(*)}$ -mesons production in e^+e^- annihilation at high energy ($M_{B_c}^2/s \gg 1$) can be described in the framework of the fragmentation model with the simple factorized expression

$$\frac{d\sigma_{B_c^{(*)}}}{dz} = \sigma_{b\bar{b}} \cdot D_{\bar{b} \rightarrow B_c^{(*)}}(z), \quad (1)$$

where $z = 2E_{B_c}/\sqrt{s}$, and $D_{\bar{b} \rightarrow B_c^{(*)}}(z)$ is the fragmentation function of $\bar{b} \rightarrow B_c$ [1-3].

The ratio of the vector B_c^* -meson production cross-section to that of the pseudoscalar B_c is

$$R = \frac{\sigma_{B_c^*}}{\sigma_{B_c}} = 1.4$$

instead of the expected value $R \simeq 3$ that one can obtain counting the quark states.

The production of the P -level states in the fragmentation model is about by an order of magnitude suppressed with respect to that for S -states [7].

The contribution of the c -quark fragmentation into the final cross-section is about two orders of magnitude lower comparing with that of the \bar{b} -quark.

ii) The production of $B_c(B_c^*)$ -mesons in the $\gamma\gamma$ collisions is described by 20 Feynmann diagrams [8-10] that can be split into three gauge invariant groups (6+6+8 diagrams, correspondingly). The first two groups describe the process of the $b\bar{b}$ and $c\bar{c}$ pairs production with the subsequent fragmentation into $B_c^{(*)}$ -meson. The contribution of the c -quark fragmentation in $\gamma\gamma$ collisions is enhanced due to the factor $(Q_c/Q_b)^4$ and can't be neglected contrary to the case of e^+e^- annihilation.

The contribution of the \bar{b} -quark fragmentation can be sufficiently well described by equation (1) and the accuracy of this description increases with the growth of p_T value. The contribution of the c -quark fragmentation is in a quite drastic contradiction with factorized expression (1). The main contribution into $B_c(B_c^*)$ -meson production at any p_T comes from the remaining 8 diagrams of the recombination type. This results in another value of $R > 3$ different from that in e^+e^- annihilation.

iii) The hadronic production of $B_c^{(*)}$ -meson requires the calculation of the full set of 36 Feynmann diagrams of the order of $O(\alpha_s^4)$ including fragmentation type diagrams. The cross section of the $B_c(B_c^*)$ can be then obtained by the convolution of the gluon-gluon cross-section with the gluon distribution functions in the initial hadrons. This leads to the dominance of the region of small $\sqrt{\hat{s}}$ of the sub-process $gg \rightarrow B_c^{(*)} + \bar{c}$ in the cross-section of B_c production. Just in the same way as in $\gamma\gamma$ collisions the recombination type diagrams dominate in gg production of $B_c(B_c^*)$ [6] and, consequently, the R value is close to 3 in the hadronic production as well.

It is obvious that the experimental cuts on p_T values rises up the role of large \hat{s} in the reaction $gg \rightarrow B_c^{(*)} + X$. In such a case we should expect the increase of the fragmentation mechanism contribution and, consequently, the simplification of all calculations. The calculations at large p_T values have been performed in [5,11] where the cross-sections of $B_c(B_c^*)$ -meson production for the FNAL Tevatron are presented. The authors of the cited papers used as the fragmentation function for $\bar{b} \rightarrow B_c^{(*)}$ the value obtained in e^+e^- annihilation for the S and P levels. Then, convoluting this fragmentation function with the cross-section of the $b\bar{b}$ pair production they calculated the distribution of $B_c(B_c^*)$ -mesons as the function of p_T . The leading logarithmic correction on the final state gluon radiation was also taken into account.

In this paper, considering the full set of Feynmann diagrams of the order of $O(\alpha_s^4)$ we will show that the fragmentation approach fails, because not all the diagrams, which can't be neglected at large p_T , were taken into account and there was some overestimation of the real phase space.

We will give a more detailed analysis of the $B_c(B_c^*)$ production at large p_T for the energy of the FNAL Tevatron ($\sqrt{s} = 1.8$ TeV).

II. HADRONIC PRODUCTION OF $B_c^{(*)}$ WITH LARGE P_T

It seems obvious at the naive level that the dominating mechanism of the $B_c^{(*)}$ -mesons production at large p_T should be connected with the fragmentation of heavy \bar{b} -quarks. The fragmentation approach is based on the assumption that the factorization of the production on the parton (\bar{b}) production with large energy and its subsequent fragmentation into different B_c states is valid. The differential cross section $d\sigma/dp_T$, for example, for $\bar{p}p$ collision will have the form

$$\begin{aligned} \frac{d\sigma}{dp_T}(\bar{p}p \rightarrow H(p_T)x) &= \sum_{i,j} \int dx_1 dx_2 dz f_{i/p}(x_1, \mu) f_{j/\bar{p}}(x_2, \mu) \times \\ &\times \frac{d\hat{\sigma}}{dp_T}(ij \rightarrow \bar{b}(p_T/z) + x) \times D_{\bar{b} \rightarrow H}(z, \mu), \end{aligned} \quad (2)$$

where $D(z, \mu)$ are fragmentation functions $\bar{b} \rightarrow H$; $H = B_c, B_c^*, \dots$, $d\hat{\sigma}/dp_T$ is $b\bar{b}$ pair production cross-section and $f_{i/A}(x, \mu)$ is parton i density function in the hadron A . This approach was used in [5,11], where the differential distributions $d\sigma/dp_T$ for $B_c^{(*)}$ -meson and its P -wave excitation were calculated. Nevertheless, it remains unclear what is the region where we can use expressions (2) and what can be the role of other subprocesses. Let us clarify these points in the framework of the same approximation – the Born-level diagrams of the order of $O(\alpha_s^4)$ [6]. The total number of diagrams in the subprocess $gg \rightarrow B_c^{(*)} + X$ is 36 and only minor part of them is connected with the \bar{b} -quark fragmentation.

We have already showed [6] that the main contribution into $B_c^{(*)}$ -meson production is connected with the diagrams of the recombination type, where both initial gluons dissociate into a pair of heavy quarks and two of these quarks recombine later into $B_c^{(*)}$ -meson. We observed the dominance of such diagrams up to the energies of gg -collisions of the order of 1 TeV. Moreover, at the energies $\sqrt{\hat{s}} \sim 30$ GeV, that give the main contribution into hadronic production of $B_c^{(*)}$ -meson it is rather meaningless to consider the fragmentation mechanism at all, because the condition $M^2/\hat{s} \ll 1$ is not valid and the pre-asymptotic terms in the definition of $D(z)$ are large. So, we can use the fragmentation approach at large values of $\sqrt{\hat{s}}$ only, but at large energies the fragmentation contribution is small as compared with that of the whole set of diagrams.

In Fig. 1 we present the results of the exact calculation for the cross-section of the subprocess $gg \rightarrow B_c(B_c^*) + X$ as a function of c.m.s. energy in comparison with that calculated in the fragmentation approach – the cross-section of the $b\bar{b}$ pair production multiplied by the probability of the fragmentation

$\bar{b} \rightarrow B_c$ and $\bar{b} \rightarrow B_c^*$, $3 \cdot 10^{-4}$ and $4.15 \cdot 10^{-4}$, correspondingly. One can see, that in the regions where the condition $\hat{s}/M_{B_c}^2 \gg 1$ is valid and expressions (2) can be used the contribution of the fragmentation is not dominant. Contrary, at small energies it gives overestimated values due to the fact that in the fragmentation calculation one uses two-particle phase space ($2 \rightarrow 2$) instead of true three-particle space ($2 \rightarrow 3$) in the partonic subprocess.

The same conclusion can be drawn from Fig. 2, where we present the differential cross-sections $d\sigma/dp_T$ for B_c - and B_c^* -mesons in comparison with the fragmentation mechanism at the energy $\sqrt{\hat{s}} = 100$ GeV in gluon-gluon collision. In the case of B_c -meson production we observe the saturation of the fragmentation at the values of $p_T > 30$ GeV, while in the case of B_c^* -meson this saturation is postponed to the very edge of the phase space ($p_T > 40$ GeV).

To stress the problems with the fragmentation picture let us consider the differential cross-section $d\sigma/dp_T$ for B_c - and B_c^* -mesons at the energy of gluon-gluon collision $\sqrt{\hat{s}} = 20$ GeV (Fig. 3). One can see that incorrect calculation of the phase space leads to the fact that the cross-section calculated by fragmentation expressions (1) exceeds that obtained from the full set of diagrams.

The conclusion one can draw from the above considerations is quite simple: at the currently reachable p_T the fragmentation approach does not work and all the diagrams contributing into $B_c^{(*)}$ mesons should be considered.

It is important to underline another point. The full and the fragmentation calculations give different predictions for the ratio of the cross-sections of B_c and B_c^* productions. The full calculation gives the value of $R \simeq 3$, while the fragmentation model predicts another value $R \simeq 1.4$.

In Fig. 4 we present the differential cross-section $d\sigma/dp_T$ for B_c and B_c^* -mesons calculated from gg cross-section convoluting it with structure functions of the initial protons for the energy of the FNAL Tevatron ($\sqrt{s} = 1.8$ TeV). We take the structure functions from [12] fixing virtuality at the value of $Q \sim 10$ GeV, because we consider only those Born diagrams contributions, which approximately correspond to those virtualities ($\alpha_s \simeq 0.2$). The running virtuality and the strong coupling constant can be used when one considers next-order radiative corrections to this process.

Two curves in the same figure correspond to the contributions calculated from expression (2) using fragmentation functions $D_{\bar{b} \rightarrow B_c}(z)$ and $D_{\bar{b} \rightarrow B_c^*}(z)$ obtained at the same order of perturbative expansion. One can see that those curves do not coincide with the exact calculations in the whole interval of p_T . At small values of p_T one can observe excess of fragmentation predictions, while at large values of p_T they are lower.

Taking into account the experimental cuts on the pseudo-rapidity ($|y| < 1$) one can calculate more realistic p_T distribution (Fig. 5). Such a cut reduces the value of cross-section by a factor of 3. It is interesting to note that in the p_T interval considered the value $R = \sigma_{B_c^*}/\sigma_{B_c}$ is about 3. Taking into account all S -wave excitations the cross-section of the B_c -meson production $1S(B_c + B_c^*) + 2S(B_c + B_c^*)$ with $p_T > 5$ GeV is about 3.3 nb that gives for Run Ib of Tevatron with the integrated luminosity $100 \div 150 \text{ pb}^{-1}$ about $3.3 \div 5.0 \cdot 10^5 B_c$ -mesons.

III. DISCUSSION

There are two different approaches to the calculation of the hadronic production of $B_c^{(*)}$ -mesons at present. The first one is based on the calculation of all the diagrams of the order of $O(\alpha_s^4)$ in the perturbative QCD. There are three publications with such calculations, which give different predictions for cross-sections [4, 6, 13]. In our previous work [6], where we considered the hadronic $B_c^{(*)}$ production, we had omitted the color factor $1/\sqrt{3}$ in the wave function of $B_c^{(*)}$ -meson that increased our prediction by a factor of 3. Correcting this point and using the same constants as in [4] we have obtained a good agreement at the level of the gg cross-sections.

The second approach is based on the usage of the fragmentation model [5,11], and, as we have demonstrated above, it gives incorrect description of the cross-section underestimating it at large energies of gg collisions and overestimating at small values of energies. So, it looks quite strange that the authors of [11] when using fragmentation approach and convoluting the resulting expressions with the structure functions have obtained the cross-sections, which are analogous to ours (see Fig. 5).

Finally, we would like to note:

- 1) In the case of the states with equal mass heavy quarks (say, ψ production) the fragmentation approach $c \rightarrow \psi$ does not also describe the exact result obtained by the calculating of all the diagrams of the order of $O(\alpha_s^4)$.

2) When one quark becomes light (u, d, s instead of c -quark) in the region of large p_T the recombination diagrams become dominant.

All these points we will consider in our next publications.

REFERENCES

1. L.Clavelli, *Phys. Rev.*, **D26**(1982), 1610;
C.-R.Ji and R.Amiri, *Phys. Rev.*, **D35**(1987), 3318;
C.-H.Chang and Y.-Q.Chen, *Phys. Lett.*, **B284**(1992), 127.
2. E.Braaten, K.Cheung and T.C.Yuan, *Phys. Rev.*, **D48**(1993), 4230.
3. V.V.Kiselev, A.K.Likhoded and M.V.Shevlyagin, *Z.Phys.*, **C63**(1994), 77.
4. C.-H.Chang and Y.-Q.Chen, *Phys. Rev.*, **D48**(1993), 4086.
5. K.Cheung, and T.C.Yuan, *Phys. Lett.*, **B325**(1994), 481.
6. A.V.Berezhnoy, A.K.Likhoded and M.V.Shevlyagin, *Preprint* IHEP-94-48 (1994), hep-ph/9408284.
7. T.C.Yuan, *Phys. Rev. Lett.*, **71**(1993), 3413.
8. A.V.Berezhnoy, A.K.Likhoded and M.V.Shevlyagin, *Phys. Lett.*, **B342**(1995), 351.
9. K.Kołodziej, A.Leike and R.Rückl, *Preprint* MPI-Ph.T/94-84, LMU-23-94 (1994).
10. F.Sartogo and M.Masetti, *private communication*.
11. K.Cheung and T.C.Yuan, *Preprint* UCD-95-4 and CPP-94- (1994).
12. J.Botts et al., “CTEQ Parton Distributions and Flavour Dependence of the Sea Quarks”, *Preprint* ISU-NP-92-17, MSUHEP-92-27 (1992).
13. S.R.Slabospitsky, *Preprint* IHEP-94-53 (1994).

FIGURE CAPTIONS

Fig.1. Gluon cross-section in nb for the B_c (white triangle) and B_c^* (black triangle) production. The fragmentation model predictions for B_c (dashed line) and B_c^* (solid line) are given for comparison.

Fig.2. $d\sigma/dp_T$ distribution in nb/GeV of $B_c(B_c^*)$ production (histograms) in gg collisions in comparison with the fragmentation approach (curves) at the energy of 100 GeV.

Fig.3. The same as in Fig.2, but for the energy of 20 GeV.

Fig.4. Differential cross-section $d\sigma/dp_T$ for $B_c(B_c^*)$ in $p\bar{p}$ -collisions at $\sqrt{s} = 1.8$ TeV in comparison with the fragmentation approach.

Fig.5. The same as in Fig.4, but with cut $|y(B_c^{(*)})| < 1$. The curves correspond to the predictions from [11].

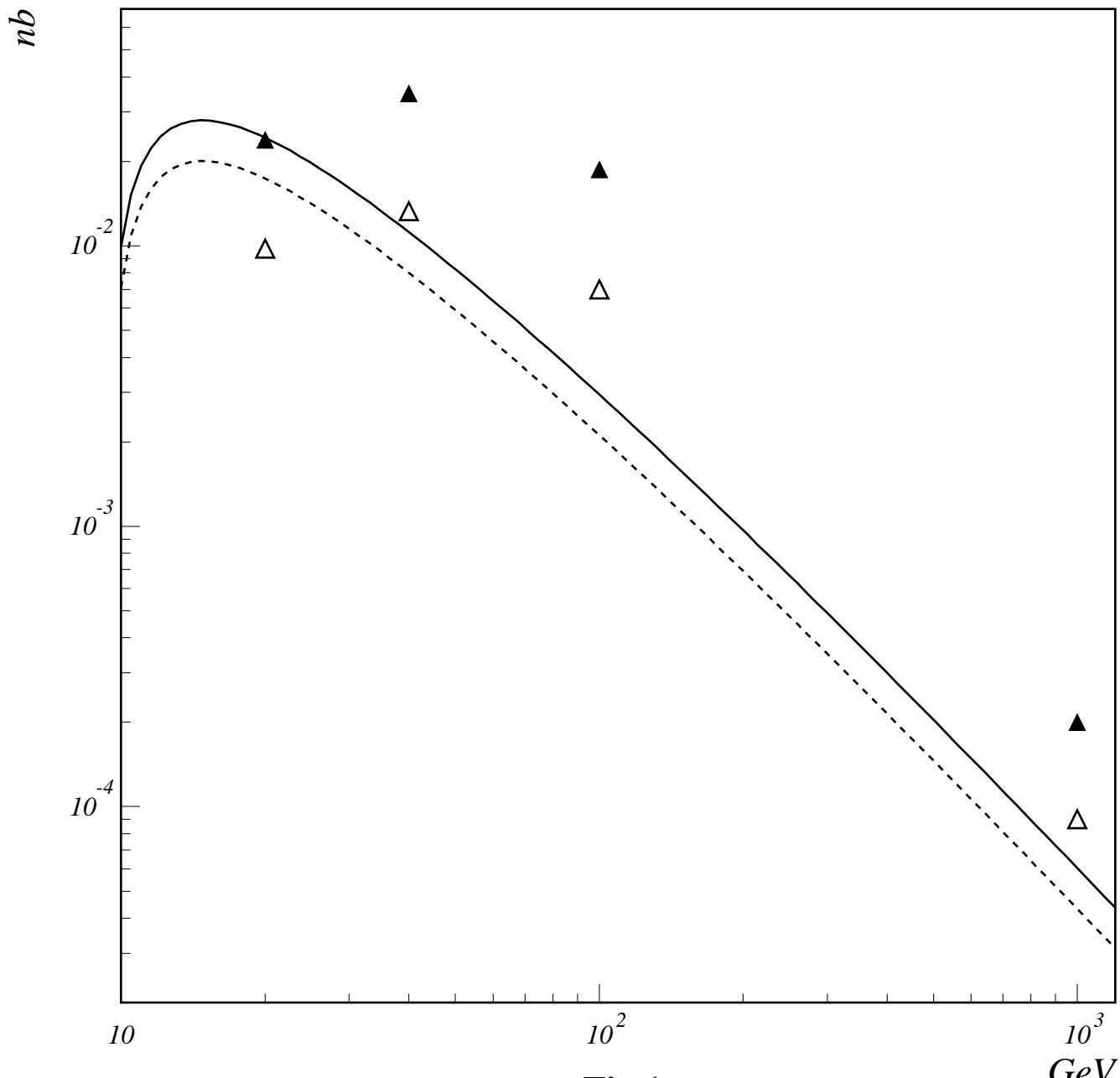


Fig.1

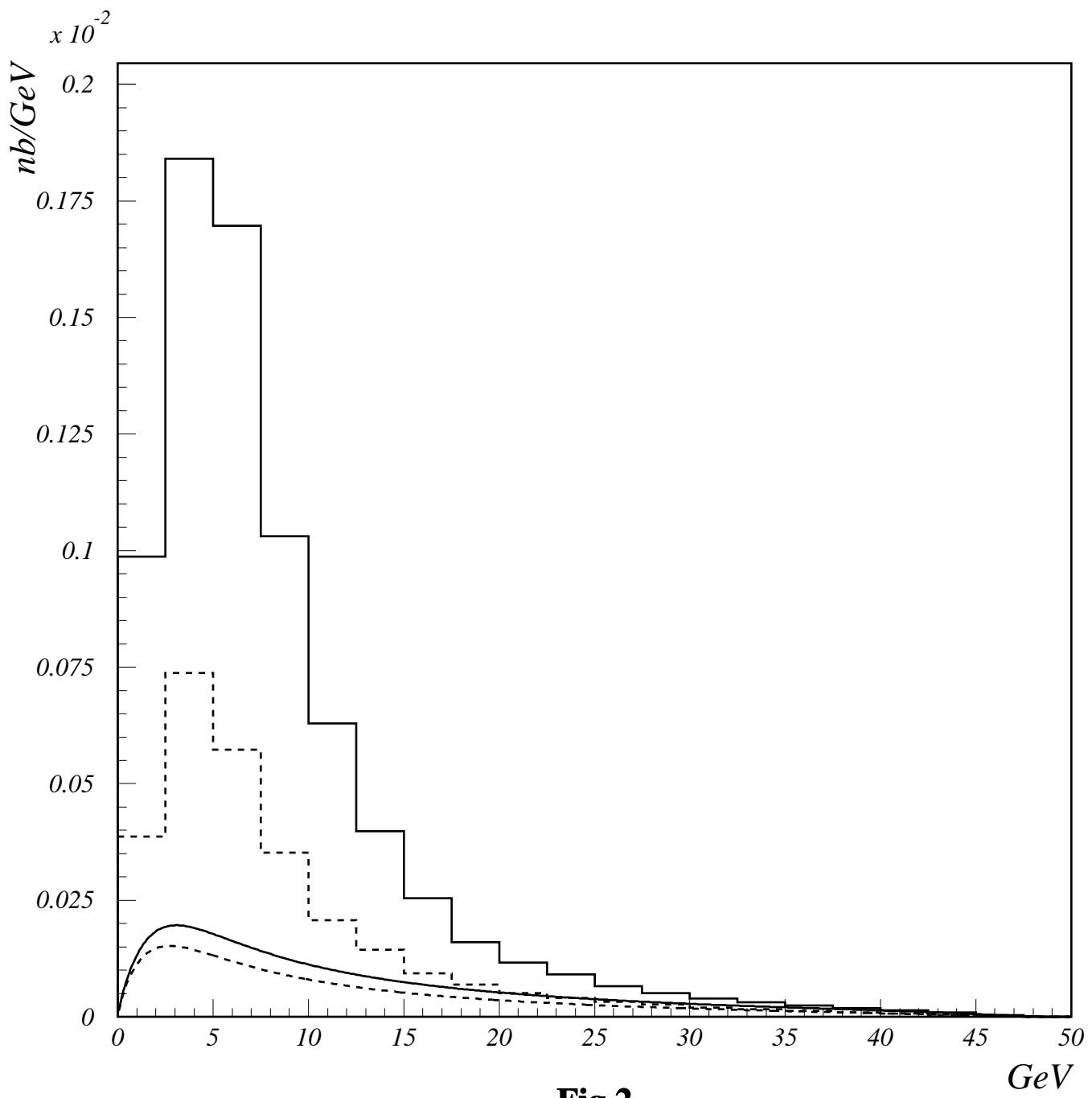
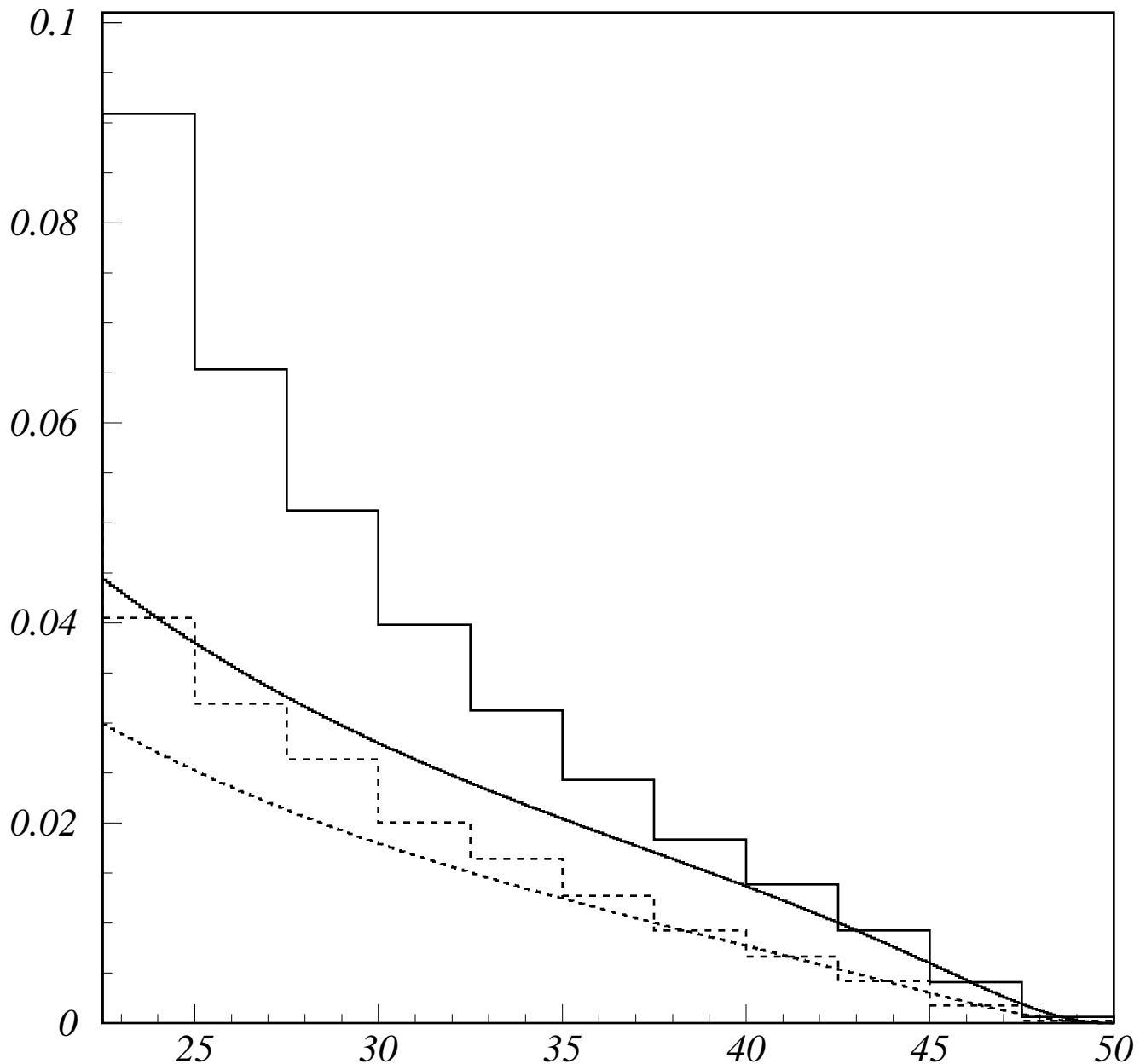


Fig.2

$x 10^{-3}$



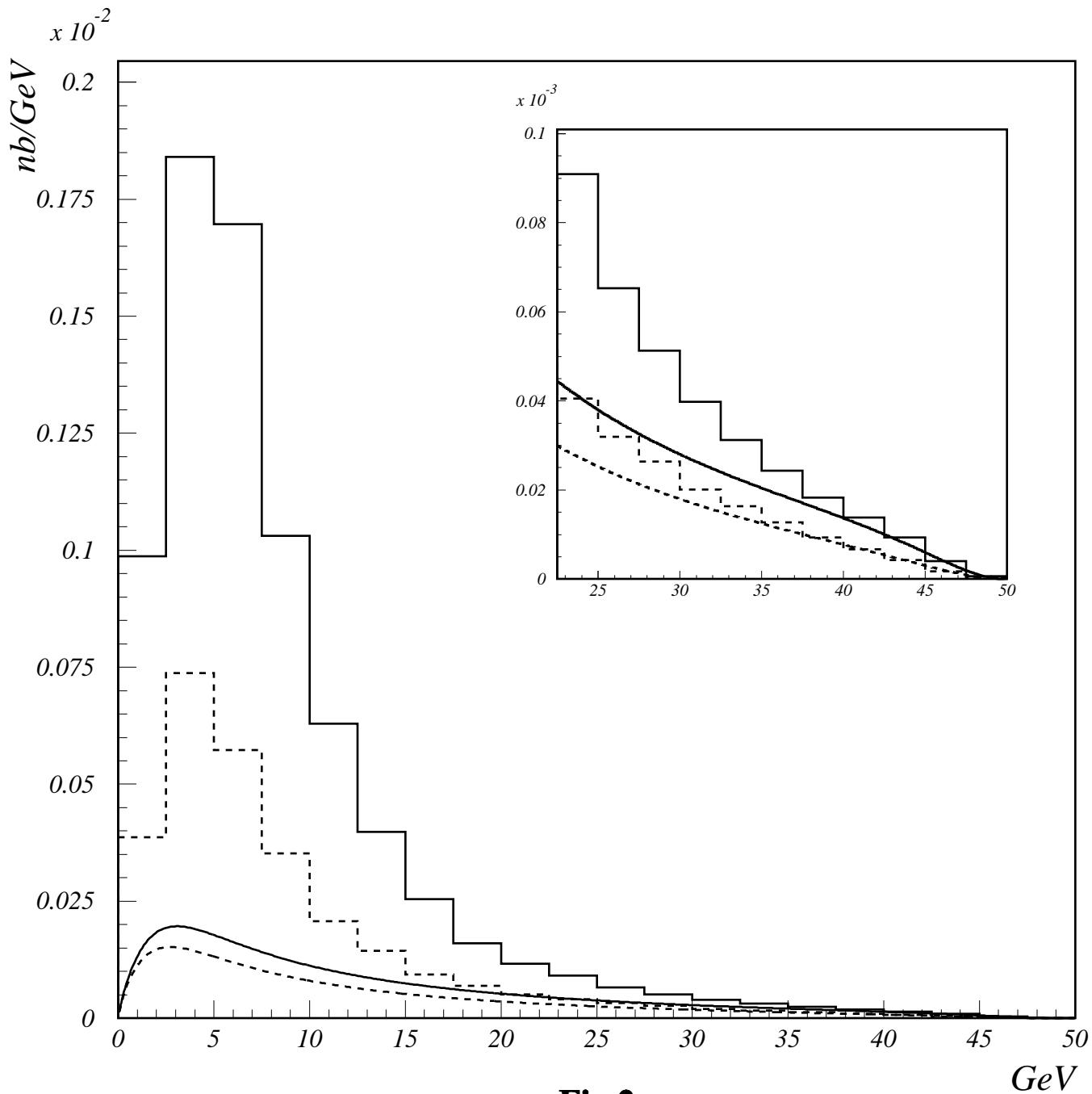


Fig.2

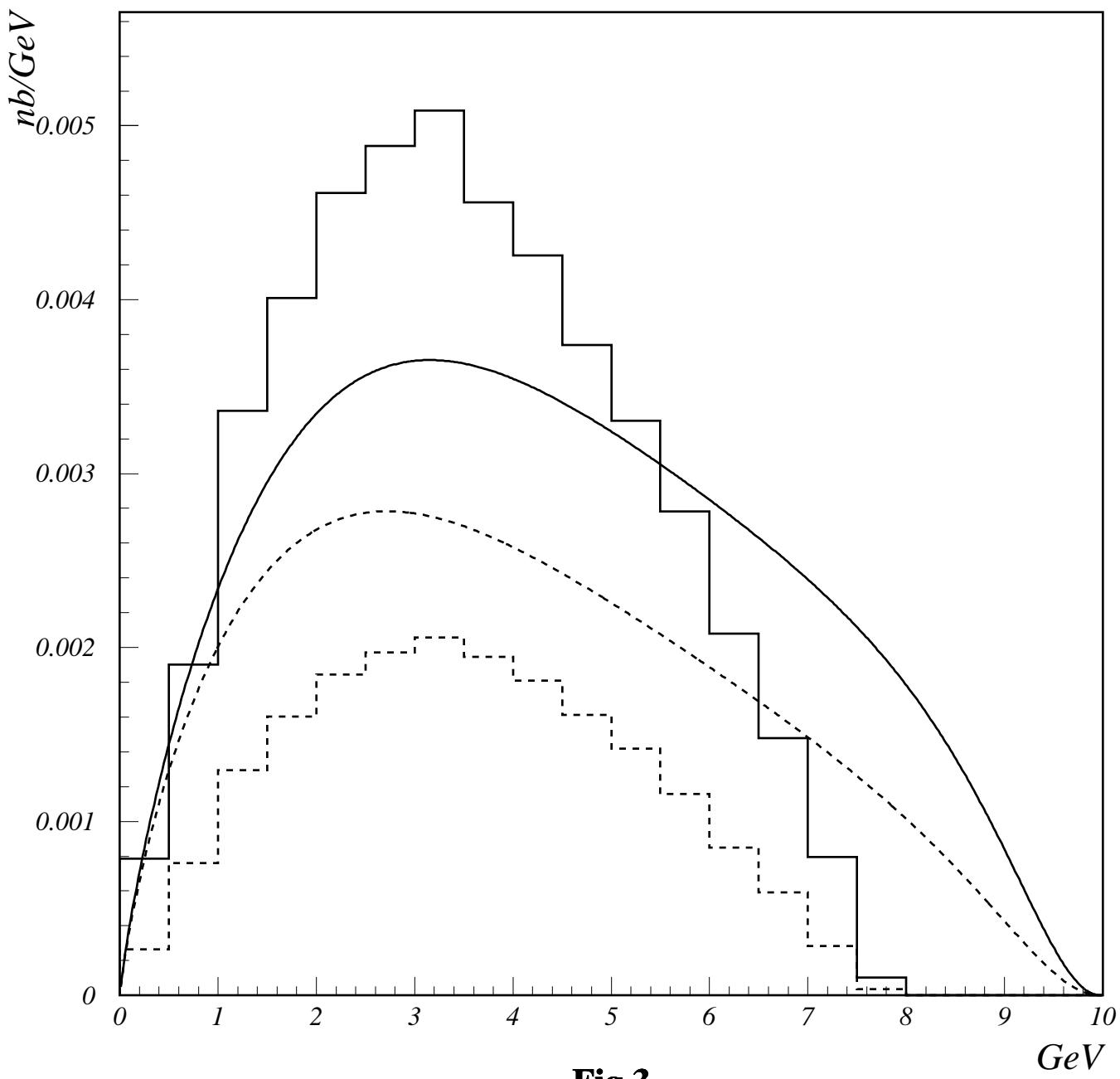


Fig.3

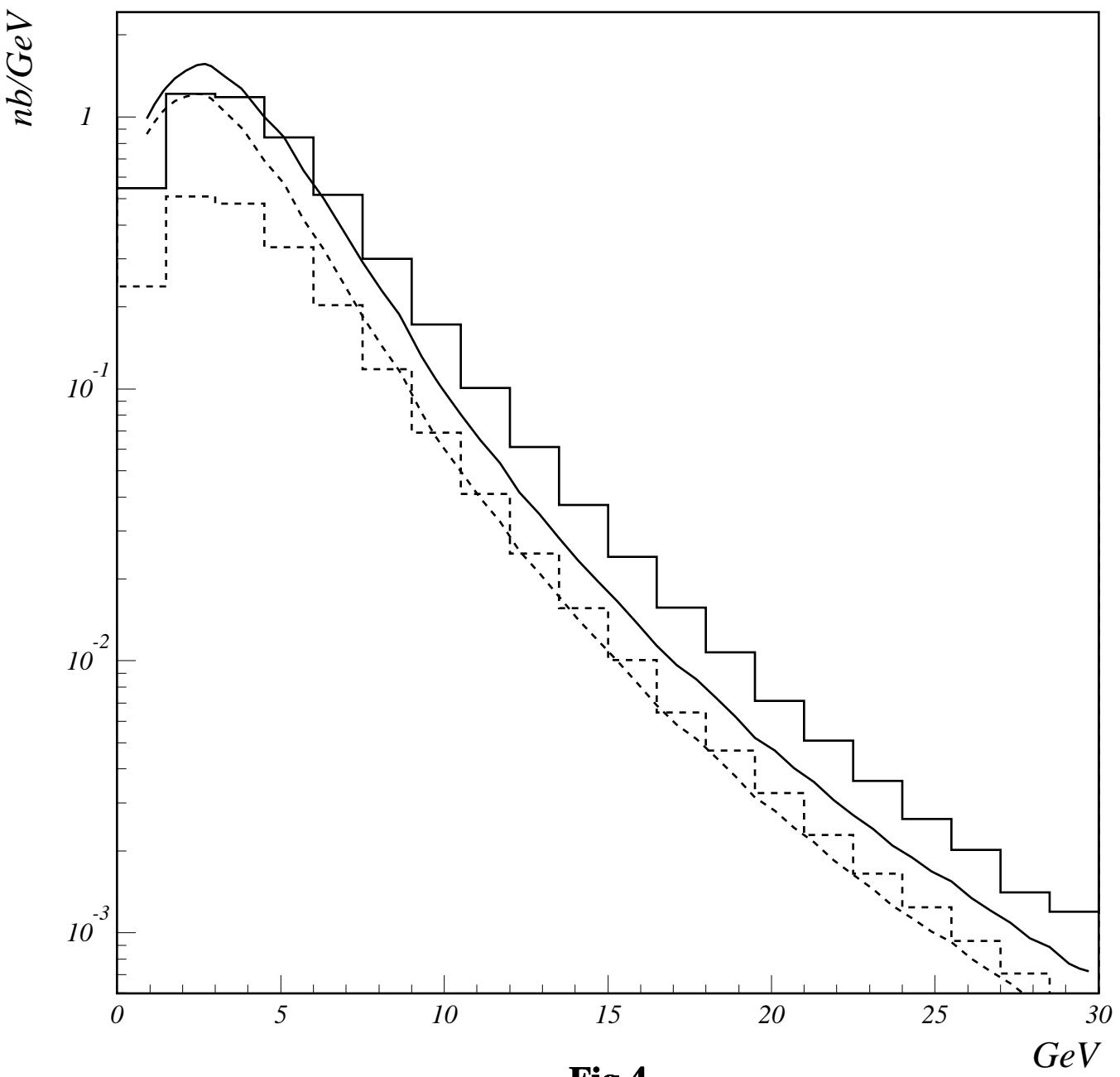


Fig.4

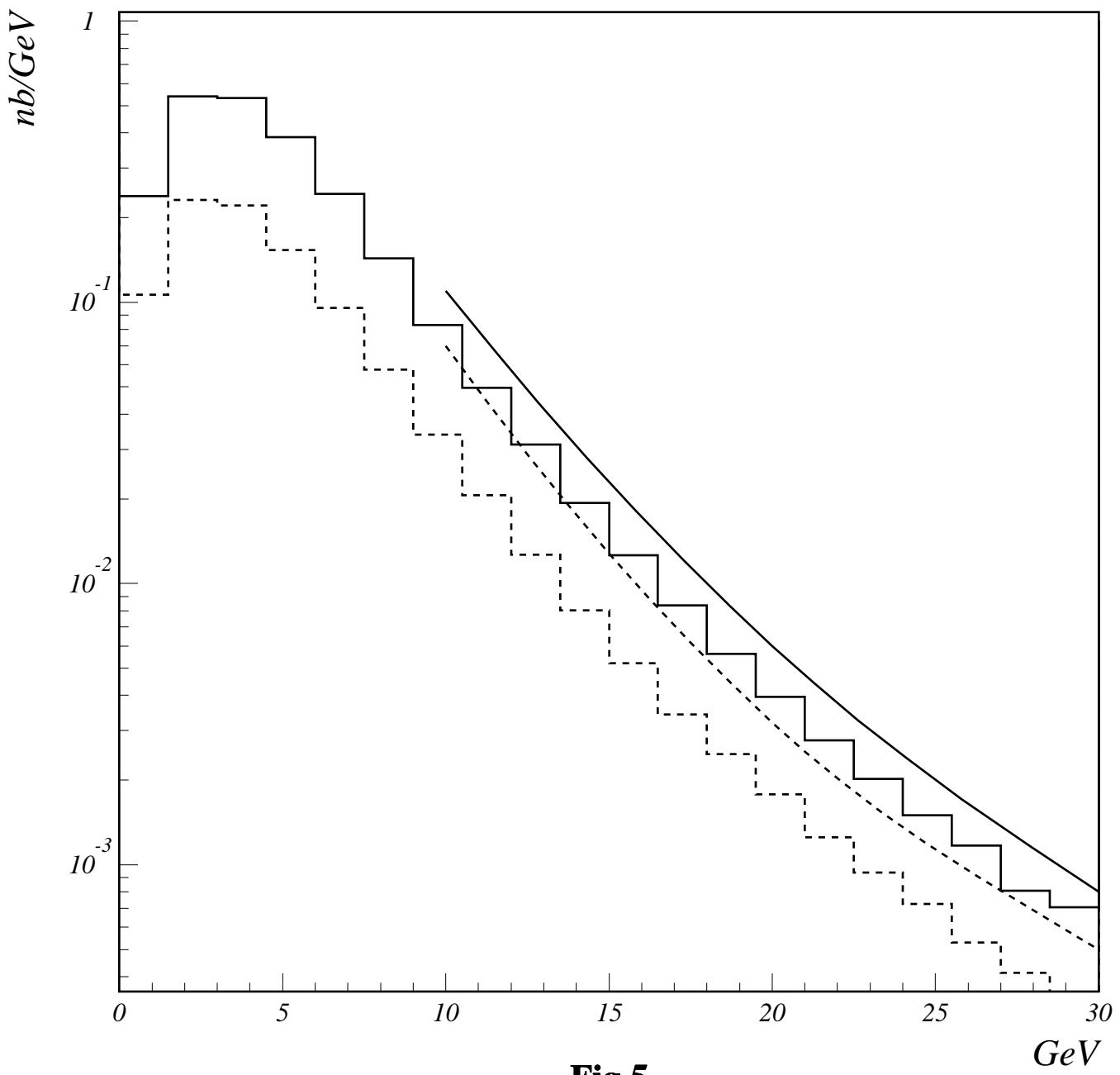


Fig.5